

Course code	MTF107
Examination date	2003-12-19
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## SOLUTIONS TO QUANTUM PHYSICS

1. i)  $Prob = |\psi_1|^2 + |\psi_2|^2 = Prob_1 + Prob_2$   
 ii)  $Prob = |\psi_1 + \psi_2|^2 = |\psi_1|^2 + |\psi_2|^2 + (\psi_1 + \psi_2)^*(\psi_1 + \psi_2) = Prob_1 + Prob_2 + 2Re(\psi_1 + \psi_2^*)$   
 The last term being due to quantum interference, giving rise to both constructive ( $2Re(\psi_1 + \psi_2^*) > 0$ ) and destructive ( $2Re(\psi_1 + \psi_2^*) < 0$ ) interference.

In i) the neutron acts as a normal, but tiny, "bullet". In ii) the neutron acts as a wave, going through both slits at once!

2. Solve

$$(S_x \cos \theta + S_y \sin \theta) \begin{pmatrix} u \\ v \end{pmatrix} = \frac{\hbar}{2} \lambda \begin{pmatrix} u \\ v \end{pmatrix}.$$

Using the  $2 \times 2$  spin matrices  $S_x, S_y$ , this gives

$$\begin{pmatrix} 0 & \cos \theta - i \sin \theta \\ \cos \theta + i \sin \theta & 0 \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix} = \lambda \begin{pmatrix} u \\ v \end{pmatrix},$$

giving

$$\begin{aligned} v e^{-i\theta} &= \lambda u \\ u e^{i\theta} &= \lambda v. \end{aligned}$$

Multiplication of left- and right-hand sides gives

$$uv(\lambda^2 - 1) = 0,$$

so the eigenvalues are

$$\lambda = \pm 1,$$

corresponding to spin-up (+1) and spin-down (-1) along the direction given by the operator.

The eigenstate to  $\lambda = +1$  satisfies

$$v = e^{i\theta} u.$$

Choosing  $u = e^{-i\theta/2}$ , gives

$$\chi_+ = N \begin{pmatrix} e^{-i\theta/2} \\ e^{i\theta/2} \end{pmatrix},$$

where  $N$  is the normalization constant, calculated as

$$\chi_+^* \chi_+ = |N|^2 (e^{i\theta/2}, e^{-i\theta/2}) \begin{pmatrix} e^{-i\theta/2} \\ e^{i\theta/2} \end{pmatrix} = 2|N|^2,$$

so

$$\chi_+ = \frac{1}{\sqrt{2}} \begin{pmatrix} e^{-i\theta/2} \\ e^{i\theta/2} \end{pmatrix}.$$

$\chi_-$  must be orthogonal to  $\chi_+$ , giving

$$\chi_- = \frac{1}{\sqrt{2}} \begin{pmatrix} e^{-i\theta/2} \\ -e^{i\theta/2} \end{pmatrix}.$$

[Note:  $\chi_+$  and  $\chi_-$  are defined only up to a constant phase, so

$$\chi_+ = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ e^{i\theta} \end{pmatrix}$$

and

$$\chi_- = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ -e^{i\theta} \end{pmatrix},$$

are also allowed, etc.]

3. The one-dimensional Schrödinger equation is

$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V(x)\psi.$$

Start from

$$\frac{\partial P}{\partial t} = \frac{\partial}{\partial t}(\psi^* \psi) = \frac{\partial \psi^*}{\partial t} \psi + \frac{\partial \psi}{\partial t} \psi^*,$$

and use the Schrödinger equation and its complex conjugate to substitute for the time derivatives in  $\frac{\partial P}{\partial t}$ , giving

$$\frac{\hbar}{2im} \frac{\partial}{\partial x} \left( \frac{\partial \psi^*}{\partial x} \psi - \psi^* \frac{\partial \psi}{\partial x} \right) = -\frac{\partial}{\partial x} \left[ \frac{\hbar}{2im} \left( \psi^* \frac{\partial \psi}{\partial x} - \psi \frac{\partial \psi^*}{\partial x} \right) \right],$$

but the quantity in the square bracket is exactly the probability current,  $j$ , so

$$\frac{\partial P}{\partial t} = -\frac{\partial j}{\partial x},$$

or

$$\frac{\partial P}{\partial t} + \frac{\partial j}{\partial x} = 0,$$

as was to be shown.

4. Generally, the time dependence of the expectation value of an operator,  $A$ , is given by

$$\frac{d}{dt} \langle A \rangle = \left\langle \frac{\partial A}{\partial t} \right\rangle + \frac{i}{\hbar} \langle [H, A] \rangle,$$

where  $A$  is an arbitrary operator, and  $H$  is the Hamiltonian operator.

For our case, this gives

$$\frac{d\langle x \rangle}{dt} = 0 + \frac{i}{\hbar} \langle [H, x] \rangle = \frac{i}{\hbar} \langle \left[ \frac{p^2}{2m}, x \right] \rangle = \frac{\langle p \rangle}{m}.$$

$$\frac{d\langle p \rangle}{dt} = 0 + \frac{i}{\hbar} \langle [H, p] \rangle = \frac{i}{\hbar} \langle -eE_0 \cos(\omega t) [x, p] \rangle = eE_0 \cos(\omega t).$$

$$\frac{d\langle H \rangle}{dt} = \left\langle \frac{\partial H}{\partial t} \right\rangle + 0 = eE_0 \omega \sin(\omega t) \langle x \rangle.$$

5. a)  $P = |\psi(x)|^2 = 4\alpha^3 x^2 e^{-2\alpha x}$  peaks when

$$0 = \frac{dP}{dx} = 4\alpha^3 2x(1 - \alpha x)e^{-2\alpha x},$$

that is, when

$$x = \frac{1}{\alpha}.$$

b)

$$\langle x \rangle = \int_0^\infty x(4\alpha^3 x^2 e^{-2\alpha x}) dx = \frac{3}{2\alpha}.$$

$$\langle x^2 \rangle = \int_0^\infty x^2(4\alpha^3 x^2 e^{-2\alpha x}) dx = \frac{3}{\alpha^2}.$$