

The Emission Spectrum due to Molecule Formation through Radiative Association

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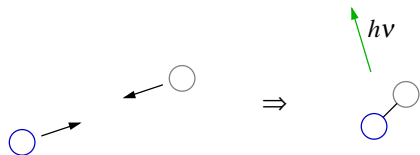
UNIVERSITY OF GOTHENBURG

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Outline

- ▶ What is radiative association?
- ▶ How is it calculated?
- ▶ Formation of HF: cross sections and emission spectra

Radiative association



- ▶ Important for molecule formation in dust-poor regions of the interstellar medium
- ▶ Stars and solar systems are born in molecular clouds, like the Orion Nebula
- ▶ Laboratory measurements difficult – only done with ionic species
- ▶ Why theoretical study of emission spectra?
 - ▶ source of radiation in planetary atmospheres
 - ▶ potential probe in laboratory

Cross section: quantum mechanical perturbation theory

$$M_{\Lambda EJ, \Lambda' v' J'} = \int_0^\infty F_{EJ}^\Lambda(R) D_{\Lambda\Lambda'}(R) \Phi_{v' J'}^{\Lambda'}(R) dR$$

E = asymptotic kinetic energy

$F_{EJ}^\Lambda(R)$ = continuum wave function

$D_{\Lambda\Lambda'}(R)$ = transition dipole between electronic states Λ and Λ'

$\Phi_{v' J'}^{\Lambda'}(R)$ = bound state wave function

$$\sigma_{\Lambda \rightarrow \Lambda' v' J'}(E) = \frac{2}{3} \frac{h^2}{(4\pi\epsilon_0)c^3} \frac{1}{2\mu E} P_\Lambda \sum_J \omega_{E\Lambda' v' J'}^3 S_{\Lambda J \rightarrow \Lambda' J'} M_{\Lambda EJ, \Lambda' v' J'}^2$$

$\omega_{E\Lambda' v' J'}$ = angular frequency of emitted photon

$S_{\Lambda J \rightarrow \Lambda' J'}$ = Hönl–London factors

$$\sigma_{\Lambda \rightarrow \Lambda'}(E) = \sum_{v' J'} \sigma_{\Lambda \rightarrow \Lambda' v' J'}(E)$$

Cross section: classical theory (no electronic transition)

Based on the Larmor power:

$$\hat{I}(b, E, \omega) = P_{\Lambda} \frac{2\omega^4}{3c^3\pi(4\pi\epsilon_0)} \left| \int_{-\infty}^{\infty} dt e^{i\omega t} \mathbf{D}_{\Lambda\Lambda}(b, E, t) \right|^2$$

with the dipole moment $\mathbf{D}_{\Lambda\Lambda}(b, E, t)$ evaluated along a trajectory with initial energy E and impact parameter b .

Spectral density:

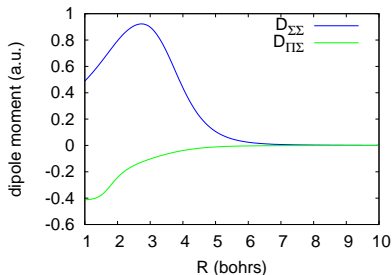
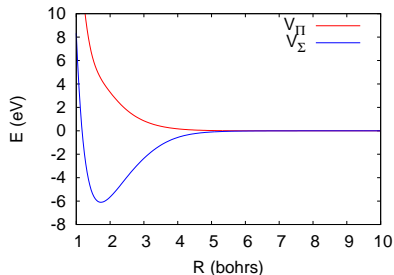
$$\frac{d\sigma_{\Lambda\rightarrow\Lambda}(E)}{d\omega} = \int_0^{\infty} 2\pi b \frac{\hat{I}(b, E, \omega)}{\hbar\omega} db$$

Radiative association cross section:

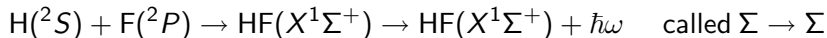
$$\sigma_{\Lambda\rightarrow\Lambda}(E) = \int_{E/\hbar}^{\omega_2} \frac{d\sigma_{\Lambda\rightarrow\Lambda}(E)}{d\omega} d\omega$$

with the limits, E/\hbar and ω_2 , defined by the range of bound states.

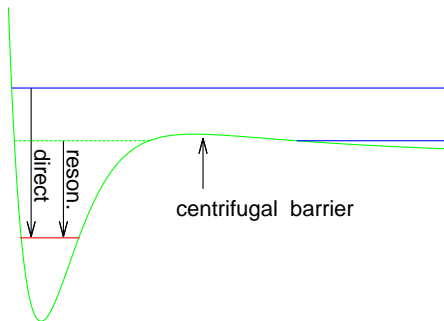
HF potentials and dipole moments



Radiative association reactions:



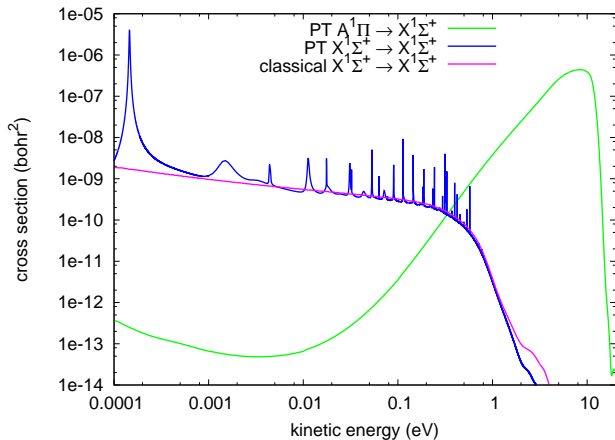
Radiative association mechanisms



HF $\Pi \rightarrow \Sigma$: only direct

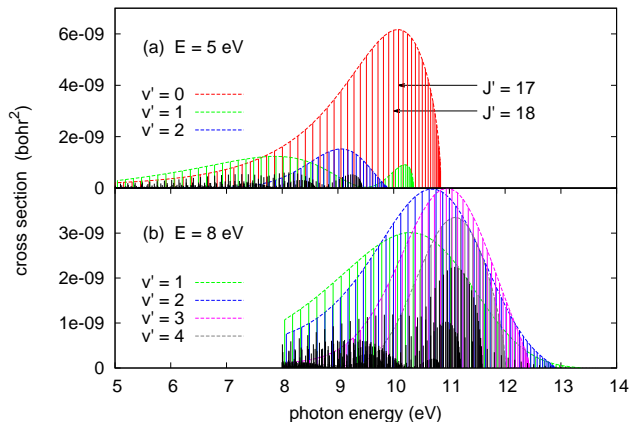
HF $\Sigma \rightarrow \Sigma$: both direct and resonance mediated

HF radiative association cross sections



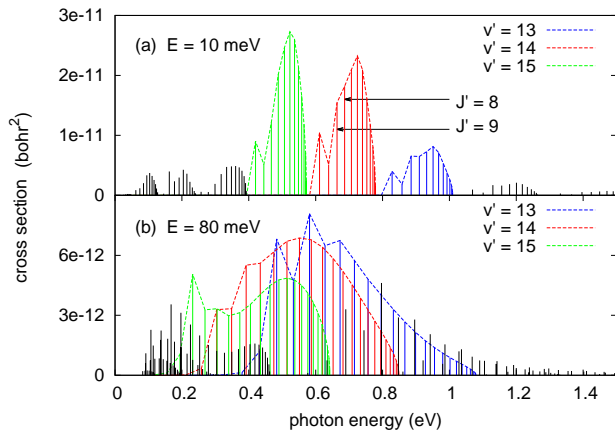
PT = quantum mechanical perturbation theory

HF $\Pi \rightarrow \Sigma$ radiative association emission spectra

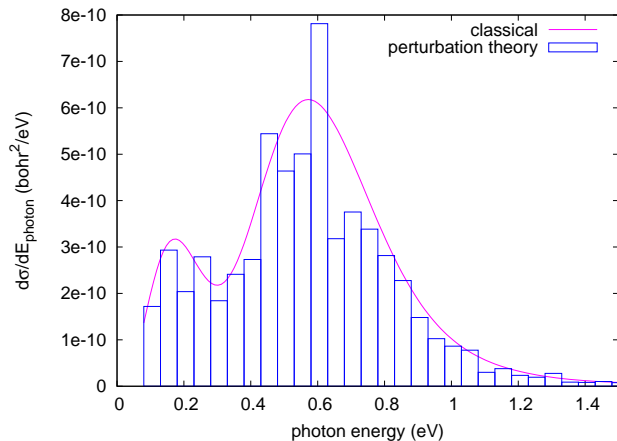


final state resolved cross section $\sigma_{\Lambda \rightarrow \Lambda' v' J'}(E)$

HF $\Sigma \rightarrow \Sigma$ radiative association emission spectra

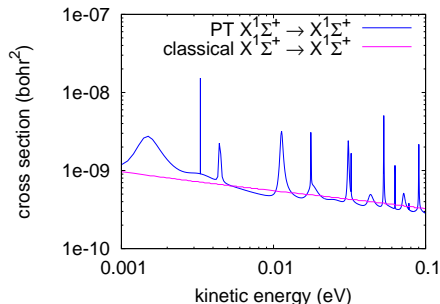


HF $\Sigma \rightarrow \Sigma$ radiative association spectral density



$E = 80 \text{ meV}$

HF $\Sigma \rightarrow \Sigma$ resonance mediated radiative association



The narrow resonance at $E = 3.304$ meV has $\nu=18$, $J=13$, and peak value $\sigma = 9.0$ bohr²

Contributions:

ν'	J'	σ (bohr ²)	E_{photon} (eV)
13	12	1.1	0.79
13	14	1.2	0.72
14	12	1.2	0.58
14	14	2.3	0.51

Conclusion

Radiative association of HF:

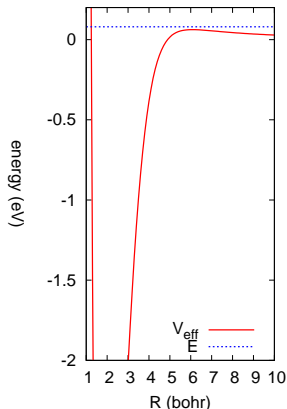
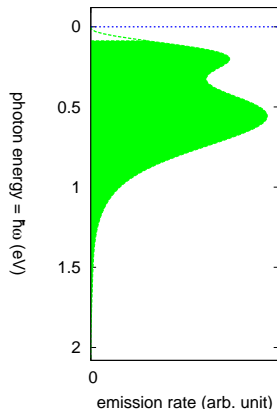
- ▶ 1–6 μm emission (IR) at low energies
- ▶ 100–150 nm emission (UV) at high energies
- ▶ resonance mediated emission could allow detection of emitted photons in laboratory experiments

Acknowledgments

- ▶ Gunnar Nyman
- ▶ The Swedish Science Council

Thank you!

Intermediate result with classical approach for H+F collisions



$E=80$ meV
 $b=6$ bohrs

emission rate $= \hat{I}(b, E, \omega)/\omega$

effective potential $= V_{\Sigma}(R) + \frac{Eb^2}{R^2}$

The shaded area under the emission curve corresponds to HF formation through radiative association.

Breit–Wigner theory

The contribution to the cross section from one resonance:

$$\sigma(E, \nu', J') = \frac{\pi \hbar^2}{2\mu E} p (2J' + 1) \frac{\Gamma_{\nu'J'}^{\text{rad}} \Gamma_{\nu'J'}^{\text{tun}}}{(E - E_{\nu'J'})^2 + (\Gamma_{\nu'J'}^{\text{tun}} + \Gamma_{\nu'J'}^{\text{rad}})^2 / 4}$$

- ▶ $E_{\nu'J'}$ is the energy of the quasi-bound level
- ▶ $\Gamma_{\nu'J'}^{\text{tun}}$ is the tunneling width
- ▶ $\Gamma_{\nu'J'}^{\text{rad}}$ is the radiative width for spontaneous emission

$E_{\nu'J'}$, $\Gamma_{\nu'J'}^{\text{tun}}$, and $\Gamma_{\nu'J'}^{\text{rad}}$ can be calculated fast without solving the full scattering problem (e.g. with LeRoy's LEVEL program).